

Light shift:

2017 P452 Q51 Q Lec 12. C.C.

Now we will quantize both atomic energy and radiation field.

atom: $\frac{\hbar}{2}\omega\sigma_z$

photons: $\hbar\omega(\hat{a}^\dagger\hat{a} + \frac{1}{2})$

atom-photon coupling $g(\sigma_+ + \sigma_-) \underbrace{(\hat{a}^\dagger + \hat{a})}_{\cos\omega t} \approx g(\sigma_+\hat{a} + \sigma_-\hat{a}^\dagger)$ } Jaynes-Cummings model

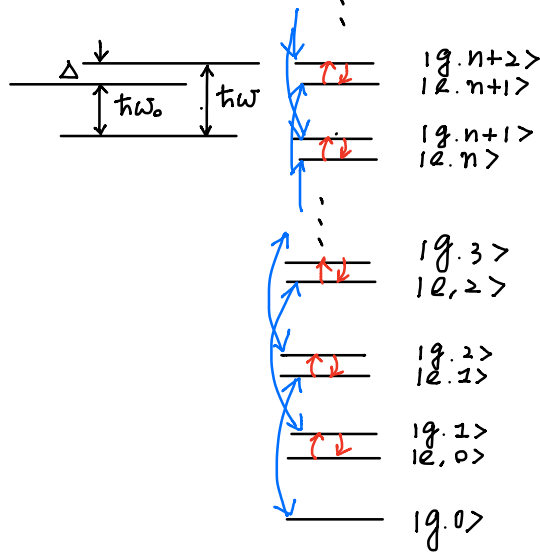
RWA

state $| \rangle = | \sigma \rangle \otimes | n \rangle$ $\sigma = 0 \text{ or } 1$ $n = 0, 1, 2, \dots$



no interaction:

assume $\Delta = \omega - \omega_0 > 0$



RW coupling $\sigma^+a + \sigma^-a^\dagger$

CRW coupling $\sigma a + \sigma^\dagger a^\dagger$

If we ignore CRW terms, all couplings occurs between $|g, n+1\rangle$ and $|e, n\rangle$

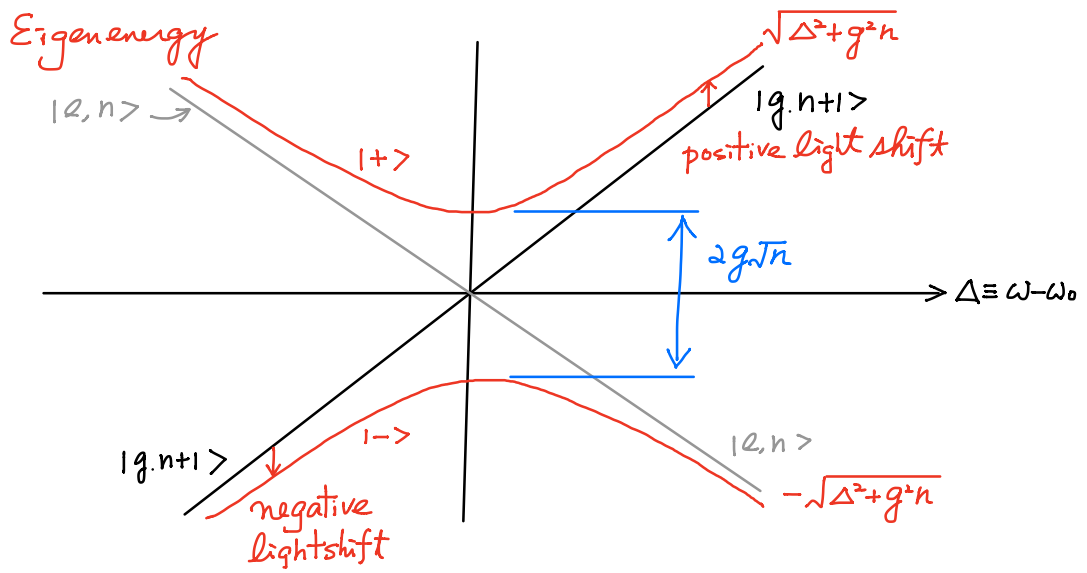
consider only a pair of coupled states: $| \rangle = \begin{pmatrix} |e, n\rangle \\ |g, n+1\rangle \end{pmatrix}$

$$H = \frac{1}{2}\hbar\omega_0\sigma_z + \hbar\omega(b^\dagger b + \frac{1}{2}) + g\sqrt{n}(\sigma + b + \sigma b^\dagger)$$

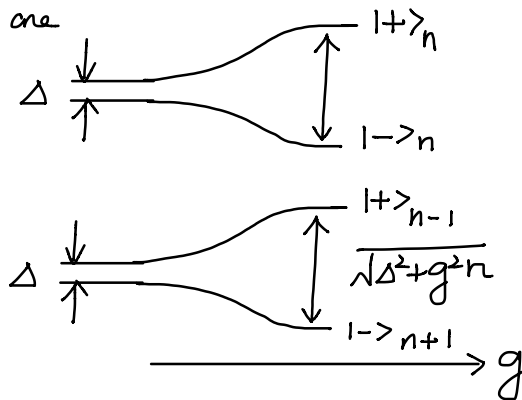
$$\langle H \rangle = \frac{1}{2}\hbar\omega_0 \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} + \hbar\omega \begin{pmatrix} n+\frac{1}{2} & 0 \\ 0 & n+\frac{3}{2} \end{pmatrix} + g\sqrt{n} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = -\frac{1}{2}\hbar\Delta\sigma_z + g\sqrt{n}\sigma_x + \text{const}\hat{1}$$

We have seen this a couple of times: Feshbach, Rabi flopping ...

This is generally how 2 modes are coupled.



Thus the eigenstates are



Thus we have a new prediction for the light shift

1. First, there are only eigenstates, $|g\rangle$ & $|e\rangle$ are "dressed" by the radiation field. This is the dressed atom picture

2. Light shift is the work done to bring an atom into the beam

$$L. \text{ shift} = \frac{1}{2} (E_+ - \Delta) \text{ for blue-detuned light } \Delta > 0$$

$$= \frac{1}{2} (E_- + \Delta) \text{ for red-detuned light } \Delta < 0$$

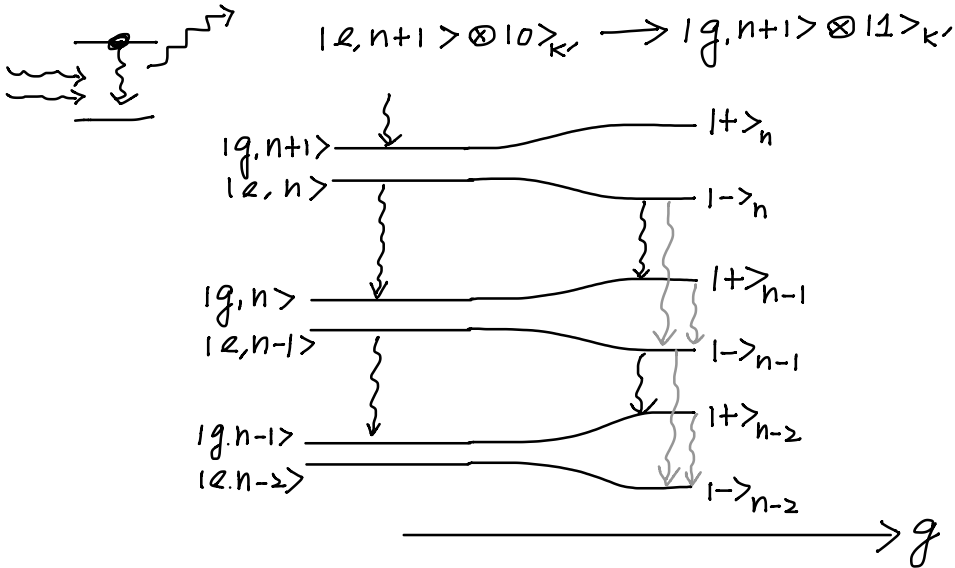
$$\Rightarrow \underline{\underline{\Delta E_L = \frac{1}{2} (\sqrt{\Delta^2 + g^2 n} - \Delta)}} \xrightarrow{\Delta \gg g\sqrt{n}} \frac{1}{4} \frac{g^2 n}{\Delta}$$

a more accurate formula!!

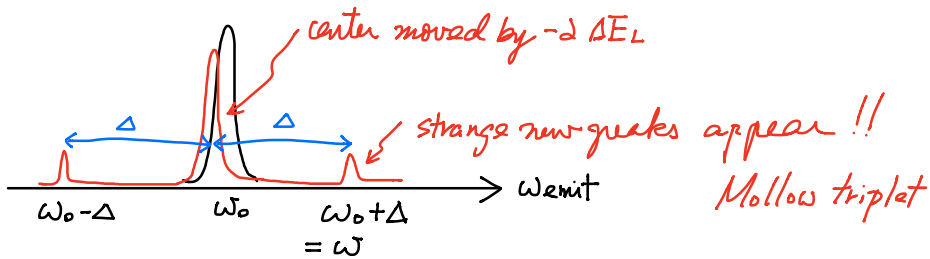
(compared to $\Delta E = \frac{P}{\hbar} \frac{I/I_s}{\Delta/P}$)

Example 1.

Can we include spontaneous emission?



Thus we predict spontaneously emitted photons should carry new components:
 $g \rightarrow 0$



Now question, illuminating some atoms with radiation at ω , we actually get emission at $\omega_0 - 2\Delta E_L \approx \omega_0$ & $\omega_0 \pm \Delta = \omega, 2\omega_0 - \omega$.

So if laser is detuned by $\Delta = -1 \text{ GHz}$, we see atoms emit photon at

$\omega_0 + 1 \text{ GHz} = \omega + 2 \text{ GHz}$. We gain 2 GHz of energy !?

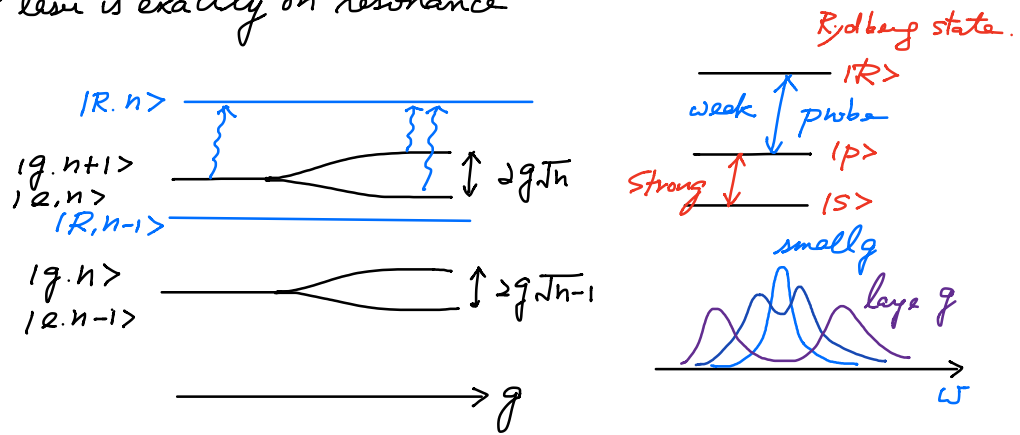
What's going on !?

Mostly, the emitted photons have the same freq as the incident photons:

Rayleigh scattering. (apart from light shift & recoil.)

Example 2: Autler-Townes effect

If laser is exactly on resonance



Many more interesting stuffs from dressed atom picture.