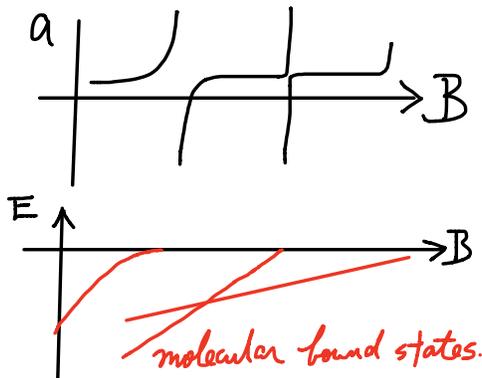


Feshbach Resonance



Features

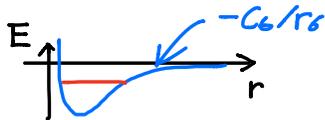
1. field dependent
2. Strong tunability
3. Associated with molecular state.
4. universal dependence

$$a = a_{bg} \left(1 - \frac{\Delta}{B - B_0} \right)$$

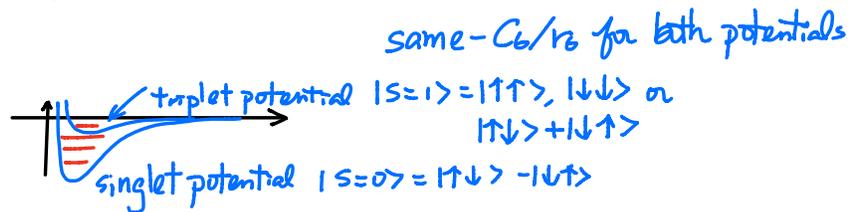
More on atomic collisions :

Molecular potential : single channel vs. multichannel scattering.

Structureless atoms: He, Ar...



Alkali-like atoms:
H, Li, Rb, Cs

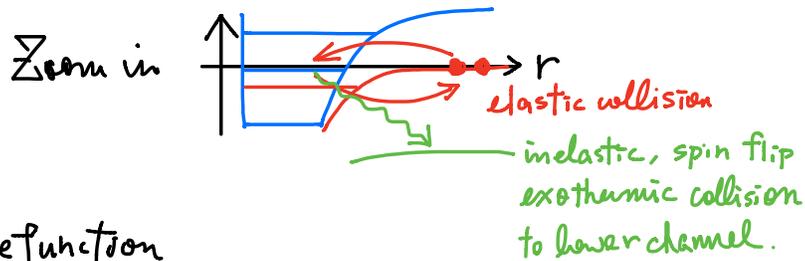
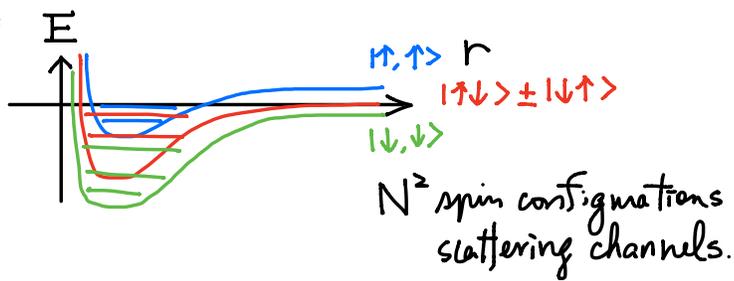


Questions: Why is singlet deeper than triplet?

More details: hyperfine structure and Zeeman shift.

$$H = \underbrace{-C_6/r^6}_{\text{Scalar shift}} + \underbrace{E_{ex}(S_1 \cdot S_2)}_{\text{exchange separate singlet \& triplet}} + \underbrace{A S_1 \cdot I_1 + A S_2 \cdot I_2}_{\text{hyperfine mix e \& nuclear spin}} - \underbrace{\mu_1 \cdot B - \mu_2 \cdot B}_{\text{Zeeman effect. magnetic field tunability.}}$$

Simplified model:



Scattering wavefunction

$$\begin{aligned} \vec{\Psi}(r) &= \frac{1}{r} e^{-ikr} \vec{\chi}_{in} - \frac{1}{r} e^{ikr} \hat{S} \vec{\chi}_{in} \\ &= \frac{1}{r} [e^{-ikr} - \hat{S} e^{ikr}] \vec{\chi}_{in} \end{aligned}$$

Examples: no collision: $\hat{S} = \hat{1} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$

only elastic collision: $S_{ij} = \delta_{ij} e^{2i\theta_i}$

cross-section $\sigma = \int |f|^2 d\Omega$, scattering ampr $f = \frac{S-1}{2ik}$
 $= \frac{4\pi}{k^2} \sin^2 \theta$

elastic collision cross-section of the i th channel: $\sigma_i = \frac{4\pi}{k^2} |S_{ii} - 1|^2$ *

inelastic collision cross section from $i \rightarrow j$: $\sigma_{ji} = \frac{4\pi}{k^2} (1 - |S_{ij}|^2)$

Probability conservation: $\sum_j |S_{ij}|^2 = 1$

For collisions between bosons, wavefunc is sym. $\Rightarrow \sigma_{el} = \frac{8\pi}{k^2} \sin^2 \theta$

if we use $k \cot \theta = -\frac{1}{a}$ at low $k \Rightarrow \sigma = \frac{8\pi a^2}{1+k^2 a^2}$

For identical Fermions, wavefunction cannot be symmetrized

$$\Rightarrow \sigma = 0.$$