

Physics 238: Atomic Physics

Fall Quarter 2021

Problem Set #6

Due: 12:00 pm, Thursday, December 2. Please submit in class.

1. Low energy excitations of a Bose-Einstein condensate

A Bose-Einstein condensate (BEC) can be described by the Gross-Pitaevskii equation under mean-field approximation,

$$i\hbar\partial_t\psi(x,t) = \left(\frac{p^2}{2m} + V(x) + g|\psi(x,t)|^2\right)\psi(x,t)$$

For a time-independent wavefunction with chemical potential μ we have

$$\left(\frac{p^2}{2m} + V(x) + g|\psi(x,t)|^2\right)\psi(x,t) = \mu\psi(x,t)$$

Here $p = -i\hbar\partial_x$ and we consider the condensate is confined in a large box $V(x) = 0$ with a uniform density n . The ground state wavefunction is thus $\psi_0(x,t) = n^{\frac{1}{2}}e^{-i\mu_0 t/\hbar}$, and the chemical potential is $\mu_0 = gn$,

- A. Here we consider low energy excited states of the system, assuming the BEC is weakly perturbed. The wavefunction can be expanded as $\psi = \psi_0 + \epsilon\psi_1$, where $\epsilon \ll 1$. Up to first order in ϵ , show that ψ_1 satisfies

$$i\hbar\partial_t\psi_1 = \left(\frac{p^2}{2m} + 2gn\right)\psi_1 + g\psi_0^2\psi_1^*,$$

which represents 2 coupled linear differential equations for ψ_1 and ψ_1^* .

- B. Apply the ansatz $\psi_1 = e^{-i\mu_0 t/\hbar}[ue^{i(kx-\omega t)} + ve^{-i(kx-\omega t)}]$, where u and v are constant amplitudes of the plane waves, and show that they satisfy

$$\begin{pmatrix} \frac{\hbar^2 k^2}{2m} + \mu_0 - \hbar\omega & \mu_0 \\ \mu_0 & \frac{\hbar^2 k^2}{2m} + \mu_0 + \hbar\omega \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix} = 0$$

- C. For solutions with non-zero amplitudes $u, v \neq 0$ show that the frequency and wave number are linked by the following dispersion $\omega(k)$

$$\omega = \frac{1}{\hbar} \sqrt{\frac{\hbar^2 k^2}{2m} \left(\frac{\hbar^2 k^2}{2m} + 2\mu_0 \right)}$$

Thus the sound speed for long wavelength excitation $k \rightarrow 0$, is $v = \lim_{k \rightarrow 0} \frac{\omega}{k} = \sqrt{\frac{\mu_0}{m}}$.

2. Low energy excitations of a Bose-Einstein condensate (second quantization)

In the second quantization form, we can write the energy of a bosonic system as

$$H = \sum_k \epsilon_k a_k^\dagger a_k + \frac{g}{2V} \sum_{k_1+k_2=k_3+k_4} a_{k_1}^\dagger a_{k_2}^\dagger a_{k_3} a_{k_4},$$

where a_k^\dagger and a_k creates and annihilates a boson with momentum k and they satisfy the bosonic commutation relation $[a_k, a_{k'}^\dagger] = \delta_{kk'}$, $\epsilon_k = \frac{\hbar^2 k^2}{2m}$ is the energy of a bare atom with momentum k and $\frac{g}{V}$ is the interaction energy of a pair of atoms.

The wavefunction of the system can be described as $|\psi\rangle = |n_1, n_2, \dots\rangle$, where n_i is the population in the i -th lowest single atom eigenstate.

- To gain some insight about the Hamiltonian, we assume there are only two momentum states $k = \pm 1$ in the system and there are only exactly 2 atoms that can occupy these states. The wavefunction can be a linear superposition of $|n_{-1}, n_1\rangle = |2, 0\rangle, |1, 1\rangle$ and $|0, 2\rangle$, where $n_{\pm 1}$ is the population in the $k = \pm 1$ state. Express the Hamiltonian as a matrix in the basis of the 3 states.
- Now we consider the system where N_0 atoms are in the lowest momentum state $|\psi_0\rangle = |N_0, 0, 0, \dots\rangle$. Show that the energy of the system is $\langle H \rangle = \frac{g}{2V} N_0(N_0 - 1) \equiv E_0$.
- Now we consider $N_0 \gg 1$ atoms in the zero momentum $k = 0$ state and few atoms $N_i \ll N_0$ in the finite momentum states. Total particle number is $N = \sum_i N_i$. Approximating $N_0 \pm 1 \approx N_0$, show that we can approximate the system energy as

$$H = E_0 + \sum_{k \neq 0} \epsilon_k a_k^\dagger a_k + \frac{g N_0}{2V} \sum_{k \neq 0} (2a_k^\dagger a_k + a_k^\dagger a_{-k}^\dagger + a_k a_{-k})$$

Remark: This result can be compared to the perturbation in 1 A.

- The Hamiltonian mixes states with k and $-k$. Show that the following Bogoliubov transformation:

$$\begin{aligned} a_k &= u_k \alpha_k + v_k \alpha_{-k}^\dagger \\ a_k^\dagger &= u_k \alpha_k^\dagger + v_k \alpha_{-k} \end{aligned}$$

Show that with suitable choice of the coefficients u_k and v_k can diagonalize the Hamiltonian as

$$H = \text{const.} + \sum_{k \neq 0} \epsilon_k \alpha_k^\dagger \alpha_k,$$

where α_k^\dagger and α_k creates and annihilates a bosonic quasi-particle with momentum k . Since the Hamiltonian is diagonal in their population $\alpha_k^\dagger \alpha_k$. These quasi-particles are effective long-lived free particles in the system, called phonons that carry sound waves and do not interact with each other.

Remark: Compare your result to 1B and 1C.