Scale invariance and universality

System remains invariant when we do the following transformation:

Simulate a giant robot walking on earth

$$\Rightarrow \begin{array}{l} m \partial_t^2 x = -mg \\ \Rightarrow x \to \lambda x \\ t \to \lambda^{1/2} t \\ \tau \to \lambda^{1/2} \tau \end{array}$$

Simulate hypothetical planet motion

$$\Rightarrow m z_4^2 r = -\frac{GMm}{r^2}$$

Motion in viscous medium

$$\Rightarrow m z_{+}^{2} \times = -mg - \varepsilon v A$$

$$X \to \lambda X$$

$$t \to \lambda^{k} t$$

$$m \to \lambda^{3} m$$

$$v \to \lambda^{k} v$$

$$mot scale invariant in least$$

$$r \rightarrow \lambda r$$
 $r \rightarrow \lambda r$
 $M \rightarrow M$ or $M \rightarrow \lambda^3 M$
 $t \rightarrow \lambda^{3/2} t$ $t \rightarrow t^\circ$
Kepler's law

not scale invariant unless $\leftarrow \Rightarrow \pi^{\frac{1}{2}} \in$

Liquid friction is 1000 times
larger than air =>
Human swimming \approx
bacteria floating in air

HW1: identify the eqn that describes water waves. How would you faithfully simulate ~100m scale ocean waves in a ~1m water tank?

Interacting particles

it
$$\partial_{t} \mathcal{G} = \mathcal{H} \mathcal{G} = \left[\sum_{i} \frac{\mathcal{P}_{i}^{2}}{2m} + V \right] \mathcal{G}$$
 $\times \rightarrow \lambda \times . \ \mathcal{P} \rightarrow \lambda^{2} \mathcal{P} . \ \mathcal{T} \rightarrow \lambda^{2} \mathcal{T} . \ V = \lambda^{-2} \mathcal{V}$

1st option: $V(x) \rightarrow V(\lambda \mathcal{V}) = \lambda^{-2} \mathcal{V}$
 $\Rightarrow V(x) = \frac{A}{x^{2}} . \ \text{Example: Etimov potential}$

and option: $V(\mathcal{P}) \rightarrow V(\lambda^{-1} \mathcal{P}) = \lambda^{-2} \mathcal{V} \Rightarrow \mathcal{V} = \mathcal{P}^{2} \quad \text{what is this?}$

Also $V = \frac{\mathcal{P}}{x} . \ \mathcal{P}^{3} \times . \ \frac{1}{x^{2}} \mathcal{Q}^{2} \mathcal{P}^{x}$

3rd: Short range interaction
$$V=g191^2$$
.

 $g'1g'1^2=g'1g1^2\lambda^{-D}=\lambda^{-2}g191^2$
 $\Rightarrow g'=g\lambda^{D-2}$
 $D=2: g=const: D gas is scale invariant.$
 $D=3: g'=g\lambda^{-1}: g=\frac{1}{x}, p. p(V=Q^2)$

Two famous examples:

2D gas with constant interaction (Stringari, 1997) Unitary Fermi gas (T.L. Ho, 2000)

HW 2: Identify two interesting scale invariant systems with interacting particles with linear energy dispersion E= p

Example: 2D gas.

itialy=
$$\int_{-\infty}^{\infty} \sqrt{1+g} |y|^2 |y|$$

This is similar to heat transport equation.

Consequence of scale invariance

$$\begin{array}{ll}
X \to \lambda X \\
P \to \lambda^{-1}P \\
E \sim P^{z} \to \lambda^{-z}E \\
+ \sim \frac{1}{E} \to \lambda^{z}E
\end{array} \Rightarrow M \to \lambda^{-D}n, \quad Y \to \lambda^{-D/2}Y$$

$$\eta = \eta(\mu, T) \Rightarrow \lambda^{D} n = n(\lambda^{2} M, \lambda^{2} T)
\Rightarrow \eta(\lambda^{2} M, \lambda^{2} T) = \lambda^{D} n(M, T)$$

if we choose
$$\lambda^z = kT \Rightarrow n(\mu.T) = \lambda^p n(\frac{M}{kT}.1)$$

 $\Rightarrow n(\mu.T) = T^{\frac{p}{2}} F(\frac{M}{kT})$

We can do this on any function

$$dE = TdS + \mu dN - pdV = 0 \Rightarrow p = Tds + \mu dn$$

$$\Rightarrow P = \int n d\mu \quad (T = const.)$$

$$= T^{D/z+1} \int_{F(x)}^{MT} dx$$

$$= T^{D/z+1} \#(\frac{M}{T}), \quad \#(x) = F(x)$$

HW3 show that entropy density s and total energy density e are given by

$$\frac{S}{N} = S = T^{\frac{1}{2}} \left[\left(\frac{D}{z} + I \right) F \left(\frac{M}{T} \right) - \frac{M}{T} F \left(\frac{M}{T} \right) \right]$$

$$\frac{E}{N} = E = T^{\frac{1}{2} + 1} \left(\frac{D}{z} \right) F \left(\frac{M}{T} \right) = n^{\frac{1+2}{2}} G \left(\frac{S}{n} \right), \text{ determine } G.$$