

Kinetic theory of gases (also developed by James Maxwell)

By mid 1800 people knew from experiments on gas

Boyle's law ($T = \text{const.}$ $P_1 V_1 = P_2 V_2$)
temperature pressure volume

Charles's law ($P = \text{const.}$ $V \propto T$)

Gay-Lussac's law ($V = \text{const.}$ $P \propto T$)

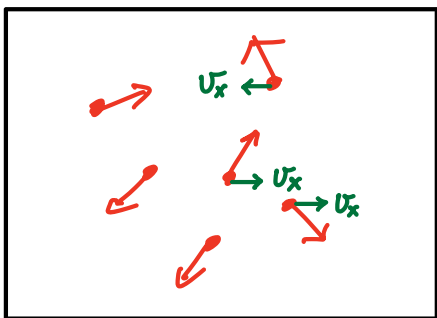
Avogadro's law ($P, T = \text{const.}$ $V \propto N$)
molecule #

$$\Rightarrow \underline{PV = N K_B T} \quad \text{or} \quad P = n K_B T$$

Boltzmann const. density = N/V

ideal gas law

Microscopic model:



Assume N molecules in a box of volume $V = L^3$.

velocity of the molecules is v .

Total energy of the system is

$$\begin{aligned} U &= N \frac{m}{2} \langle v^2 \rangle = N \frac{m}{2} (\langle u_x^2 \rangle + \langle u_y^2 \rangle + \langle u_z^2 \rangle) \\ &= \frac{3}{2} N m \langle u_x^2 \rangle \\ &\equiv N \epsilon_K \quad \text{mean kinetic energy per molecule.} \end{aligned}$$

Pressure on the wall = Force/area

$$P = \frac{F}{A} = \frac{\langle \Delta P / \Delta t \rangle}{L^2} = \frac{1}{L^2} N \langle \frac{2m u_x}{2L / u_x} \rangle = \frac{N}{L^2} \frac{m \langle u_x^2 \rangle}{L} = \frac{N}{V} \frac{m}{3} u^2 = \frac{2}{3} n \epsilon_K$$

$$\text{Compare with } P = n K_B T \Rightarrow \epsilon_K = \frac{3}{2} K_B T \quad \text{or} \quad \frac{m}{2} \langle u_x^2 \rangle = \frac{m}{2} \langle u_y^2 \rangle = \frac{m}{2} \langle u_z^2 \rangle = \frac{1}{2} K_B T$$

$$\Rightarrow U = N \epsilon_K$$

Equipartition theorem: in equilibrium, each degree of freedom has a mean energy $\frac{1}{2} K_B T$.

Atomic gas: $\epsilon_k = \frac{3}{2}KT$ (translational energy) $\Rightarrow$ rms velocity $\sqrt{\langle v^2 \rangle} = \sqrt{\frac{3KT}{m}}$

molecular gas: $\epsilon_k = \frac{3}{2}KT + KT + \text{vibrational} + \text{electronic} + \text{nuclear} \dots$
at room temperature. $\nwarrow$ rotational energy

$$\Rightarrow U = N\gamma KT \quad \gamma = \frac{3}{2} \text{ (atoms)}, \frac{5}{2} \text{ (regular molecules)}$$

What's the probability distribution $P(v_x, v_y, v_z)$. $\int P(\vec{v}) d^3v = 1$?

Statistical mechanics: probability $P(E)$ of a particle with energy E is

$$P(E) = \frac{1}{Z} e^{-E/KT}, \text{ where } Z \text{ is the normalization const.}$$
$$\sum_E \frac{1}{Z} e^{-E/KT} = 1 \Rightarrow Z = \sum_E e^{-E/KT}.$$

Maxwell's approach: Probability $P(v_x, v_y, v_z)$ should be independent in all directions

$$P(v) = p(v_x)p(v_y)p(v_z) \quad v = \sqrt{v_x^2 + v_y^2 + v_z^2}$$

$$\partial_{v_x} P(v) = P'(v) \frac{dv}{dv_x} = P'(v) \frac{v_x}{v} = p'(v_x)p(v_y)p(v_z)$$

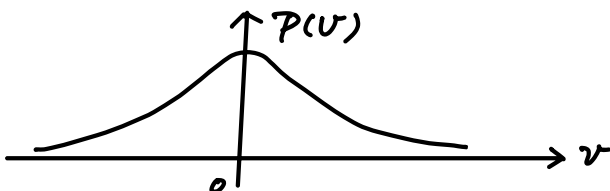
$$\Rightarrow \frac{P'}{P} = \frac{p'}{p} = \text{const} \Rightarrow \frac{dP}{P} = C v dv \quad \frac{dp}{p} = C v_x dv_x$$

$$\Rightarrow P(v) = C_1 e^{-Cv^2}, \quad p(v_x) = C_2 e^{-Cv_x^2}$$

$$\Rightarrow P(v) = C_1 e^{-\beta v^2} = C_1 e^{-\beta v_x^2} e^{-\beta v_y^2} e^{-\beta v_z^2}$$

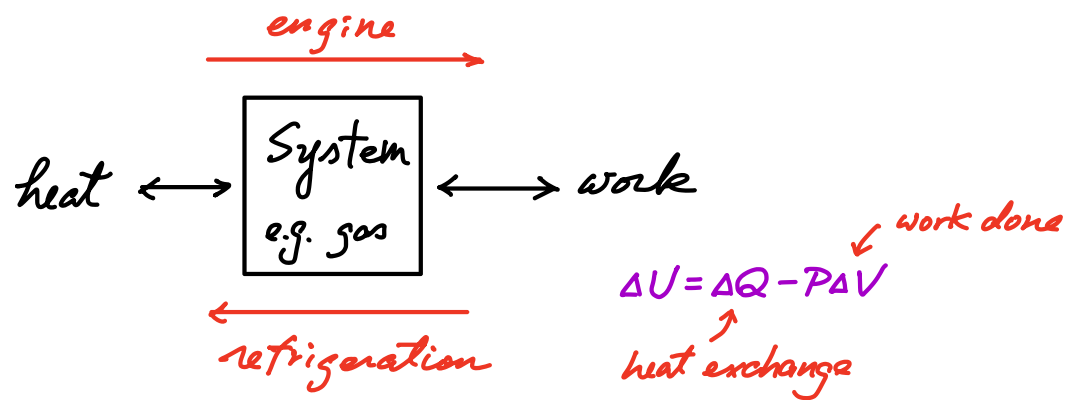
Using $\int P(v) dv = 1$ and $\langle \frac{1}{2}mv^2 \rangle = \frac{3}{2}KT$. We get

$$P(v_x, v_y, v_z) = \left(\frac{m}{2\pi KT}\right)^{3/2} e^{-\beta \frac{m}{2}(v_x^2 + v_y^2 + v_z^2)}, \quad \beta = 1/KT$$



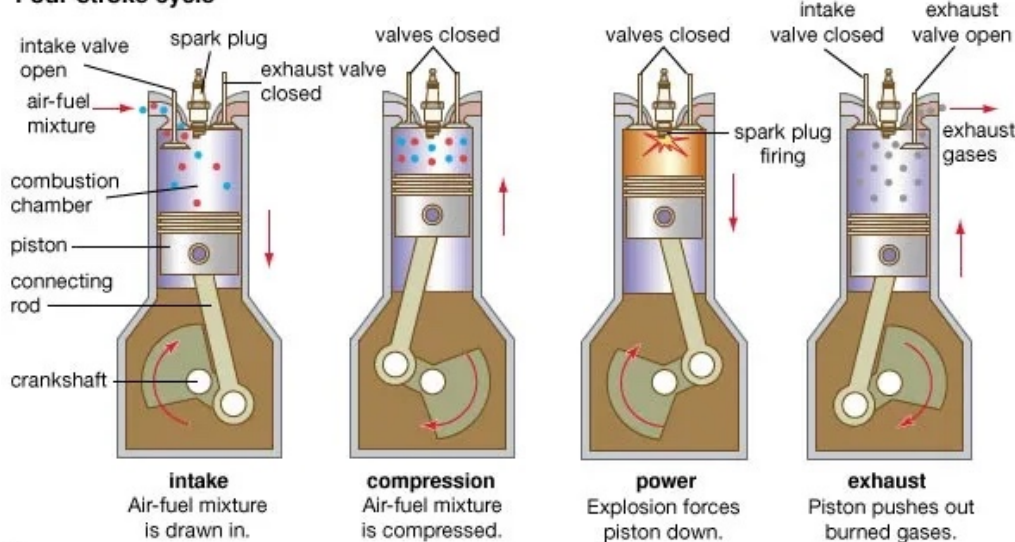
Heat engine

Thermodynamics was developed to understand and improve heat engine: Internal combustion engine (ICE), refrigeration.



We will use engine to illustrate how engine does work.

Four-stroke cycle



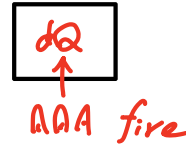
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Intake stroke → combustion stroke
→ Ignition stroke
→ Power stroke
→ exhaust stroke

Engine absorbs heat and converts it into work.

Thermodynamical processes:

1. Isochoric process ($V = \text{const.}$)



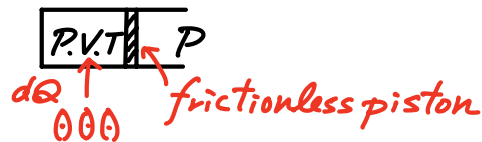
heat transfer $\rightarrow dQ = C_v M dT$ $\leftarrow$ temperature increase

isochoric specific heat $\uparrow$ mass $\uparrow$

$$dQ = dU = d[N \gamma k T] \Rightarrow C_v = \frac{1}{M} \frac{dU}{dT} = \frac{\gamma k}{m}$$

$\uparrow$
molecular mass

2. Isobaric process ($P = \text{const.}$)



$$dU = dQ - F dx = dQ - P dV$$

isobaric specific heat $C_p \equiv \frac{1}{M} \frac{dQ}{dT} = \frac{1}{M} \left(\frac{dU}{dT} + P \frac{dV}{dT} \right)$ $\leftarrow PV = NkT$

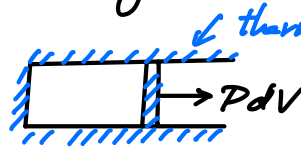
$$\Rightarrow \frac{dV}{dT} \Big|_{P=c} = \frac{Nk}{P}$$

$$= \frac{N}{M} \gamma k + \frac{N}{M} k = \frac{N}{M} (\gamma + 1) k = \frac{\gamma + 1}{\gamma} C_v$$

3. Adiabatic process (no heat exchange, isolated system)

$$dU = -P dV = \gamma N k dT$$

$$= \gamma N k d(PV)$$



$$\Rightarrow P dV + \gamma P dV + \gamma V dp = 0$$

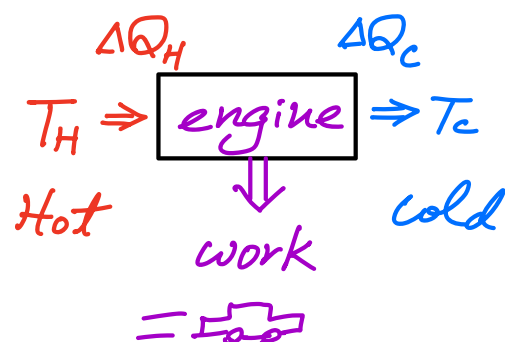
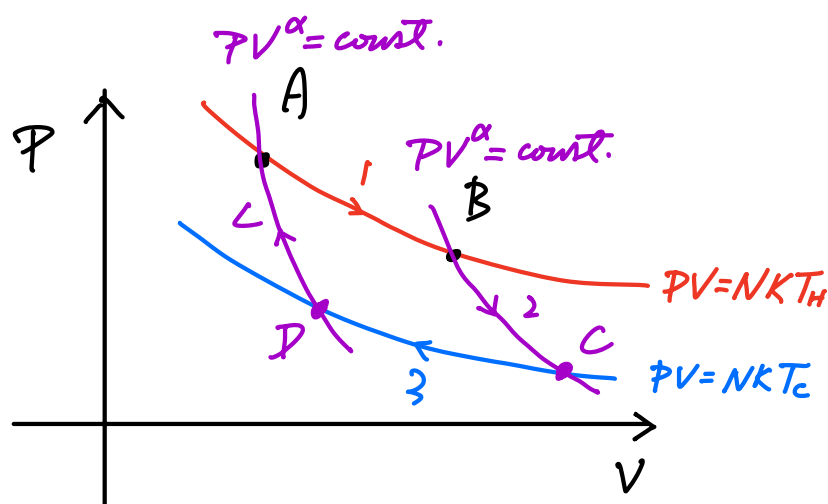
$$\Rightarrow (\gamma + 1) P dV + \gamma V dp = 0 \Rightarrow (\gamma + 1) \frac{dV}{V} + \gamma \frac{dP}{P} = 0 \Rightarrow PV^{\frac{\gamma+1}{\gamma}} = \text{const.}$$

4. Isothermal process ($T = \text{const.}$) $\Rightarrow PV = c$. $U = c$.

$$\Rightarrow dU = dQ - P dV = 0 \Rightarrow dQ = P dV$$

Carnot Cycle: The most efficient process we know for a heat engine.

System absorbs heat from T_H , does work, and cools down by T_C .



step 1: isothermal expansion ($T = T_H$, $V_A \rightarrow V_B$)

$$\Delta U = \Delta Q - P\Delta V = 0. \quad \Delta Q = P\Delta V = \int_A^B P dV$$

Absorb energy from heat source and do work by expanding

step 2: adiabatic expansion ($T_H \rightarrow T_C$, $\Delta Q = 0$)

Cool down the system and do some more work.

step 3: isothermal contraction ($T = T_H$, $V_C \rightarrow V_D$)

Contract and dump excess heat to cold reservoir.

step 4: adiabatic contraction ($T_C \rightarrow T_H$, $\Delta Q = 0$)

More contraction and heat system back to T_H

Energy absorbed from the hot source @ T_H

$$\Delta Q = P \Delta V = \int_A^B P dV = \int_A^B \frac{NKT_H}{V} dV = NKT_H \ln \frac{V_B}{V_A}$$

Energy released into the cold source @ T_C

$$\Delta Q = P \Delta V = \int_C^D P dV = \int_C^D \frac{NKT_C}{V} dV = NKT_C \ln \frac{V_D}{V_C}$$

During steps 2 and 4, no heat exchange we only need to evaluate work $\int P dV$

$$\text{Step 2 Work done} = \int_B^C P dV = P_B V_B^\alpha \int_B^C V^{-\alpha} dV = P_B V_B^\alpha \frac{1}{-\alpha+1} (V_C^{-\alpha+1} - V_B^{-\alpha+1})$$

$$\text{Step 4 Work done} = P_A V_A^\alpha \frac{1}{-\alpha+1} (V_A^{-\alpha+1} - V_D^{-\alpha+1})$$

Total work done by the system:

$$ABCD \text{ area} = \int_A^B + \int_B^C + \int_C^D + \int_D^A P dV$$

$$= NK (T_H \ln \frac{V_B}{V_A} + T_C \ln \frac{V_D}{V_C}) + NK T_H \frac{1}{1-\alpha} \left[\left(\frac{V_B}{V_C} \right)^{\alpha-1} - 1 + 1 - \left(\frac{V_A}{V_D} \right)^{\alpha-1} \right]$$

$$\text{Step 1: } P_A V_A = P_B V_B \quad \text{Step 2: } P_B V_B^\alpha = P_C V_C^\alpha \quad \Rightarrow \quad \frac{V_B^\alpha}{V_C^\alpha} = \frac{P_C}{P_B} = \frac{P_D}{P_A} \frac{V_D}{V_C} \frac{V_B}{V_A} = \frac{V_A^\alpha V_D V_B}{V_D^\alpha V_A V_C}$$

$$\text{Step 3: } P_C V_C = P_D V_D \quad \text{Step 4: } P_D V_D^\alpha = P_A V_A^\alpha \quad \Rightarrow \quad V_A/V_D = V_B/V_C$$

$$\text{Step 2 \& 4 do zero net work: } \int_B^C + \int_D^A P dV = 0.$$

$$\Rightarrow ABCD \text{ area} = NK (T_H - T_C) \ln \frac{V_B}{V_A}$$

$$\text{Carnot cycle efficiency} = \frac{\text{work}}{\text{energy needed}} = \frac{NK (T_H - T_C) \ln \frac{V_B}{V_A}}{NK T_H \ln \frac{V_B}{V_A}} = 1 - \frac{T_C}{T_H}$$